

# Comparing Trimaximal Mixing and Its Variants with Deviations from Tri-bimaximal Mixing

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## Abstract

We analyze in detail the predictions of “trimaximal” neutrino mixing, which is defined by a mixing matrix with identical second column elements. This column is therefore identical to the second column in the case of tri-bimaximal mixing. We also generalize trimaximal mixing by assuming that the other rows and columns of the mixing matrix individually can have the same forms as for tri-bimaximal mixing. The phenomenology of these alternative scenarios and their mixing angle and CP phase correlations are studied. We emphasize how trimaximal mixing scenarios can be distinguished experimentally from broken tri-bimaximal mixing.

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# 1 Introduction

Ten years after observation of the depletion of atmospheric muon-neutrinos was established by the SuperKamiokande collaboration [1], our knowledge of the neutrino oscillation parameters has been noticeably sharpened by the ensuing atmospheric [2], solar [3], reactor [4], and long-baseline [5] neutrino experiments. Recent global analyses [6–8] limit the oscillation parameters to the  $1\sigma$  and  $3\sigma$  ranges determined by Fogli *et al.*, for example, to be [8]

$$\begin{aligned}\Delta m_{32}^2 &= \left(2.39_{-0.08, 0.33}^{+0.11, 0.42}\right) \times 10^{-3} \text{ eV}^2, \\ \Delta m_{21}^2 &= \left(7.67_{-0.19, 0.53}^{+0.16, 0.52}\right) \times 10^{-5} \text{ eV}^2, \\ \sin^2 \theta_{23} &= 0.466_{-0.058, 0.135}^{+0.073, 0.178}, \\ \sin^2 \theta_{12} &= 0.312_{-0.018, 0.049}^{+0.019, 0.063},\end{aligned}\tag{1}$$

while the reactor neutrino oscillation angle remains more uncertain. In a recent analysis [8] it is weakly ( $1.6\sigma$ ) constrained to be non-zero according to (see also Ref. [9])

$$\sin^2 \theta_{13} = 0.016 \pm 0.010 \ (\leq 0.046).\tag{2}$$

These mixing angles help to specify the Pontecorvo-Maki-Nakagawa-Sakata (PMNS) mixing matrix defined in the standard convention [10] by

$$U_{\text{PMNS}} = \begin{pmatrix} c_{12} c_{13} & s_{12} c_{13} & s_{13} e^{-i\delta} \\ -s_{12} c_{23} - c_{12} s_{23} s_{13} e^{i\delta} & c_{12} c_{23} - s_{12} s_{23} s_{13} e^{i\delta} & s_{23} c_{13} \\ s_{12} s_{23} - c_{12} c_{23} s_{13} e^{i\delta} & -c_{12} s_{23} - s_{12} c_{23} s_{13} e^{i\delta} & c_{23} c_{13} \end{pmatrix},\tag{3}$$

where  $c_{ij} = \cos \theta_{ij}$ ,  $s_{ij} = \sin \theta_{ij}$  with  $\delta$  the unknown CP-violating Dirac phase. Harrison, Perkins, and Scott first emphasized that the experimentally obtained mixing matrix is close to the simple tri-bimaximal mixing (TBM) form where [11]

$$U_{\text{TBM}} = \begin{pmatrix} \sqrt{\frac{2}{3}} & \sqrt{\frac{1}{3}} & 0 \\ -\sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & -\sqrt{\frac{1}{2}} \\ -\sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & \sqrt{\frac{1}{2}} \end{pmatrix}.\tag{4}$$

The column vectors are just the eigenvectors of the three diagonal  $U(3)$  operators. For exact tri-bimaximal mixing, the mixing angles are<sup>1</sup>

$$\sin^2 \theta_{12} = \frac{1}{3}, \quad \sin^2 \theta_{23} = \frac{1}{2}, \quad \sin^2 \theta_{13} = 0,\tag{5}$$

which are seen to be close to the present data quoted above in Eq. (1).

In a top-down approach, the charged lepton mass matrix,  $m_\ell$ , and light Majorana neutrino mass matrix,  $m_\nu$ , are specified in some model with particular family and flavor symmetries, often by invoking the seesaw mechanism. The two mass matrices are diagonalized by two unitary transformations such that

$$\begin{aligned}U_\ell^\dagger m_\ell^\dagger m_\ell U_\ell &= \text{diag}(m_e^2, m_\mu^2, m_\tau^2), \\ U_\nu^T m_\nu U_\nu &= m_\nu^{\text{diag}} = \text{diag}(m_1, m_2, m_3).\end{aligned}\tag{6}$$

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<sup>1</sup>The experimental results are so close to TBM that parametrizations of the PMNS matrix with TBM as the starting point have been proposed [12].

The PMNS mixing matrix then follows from  $U_{\text{PMNS}} P = U_\ell^\dagger U_\nu$ , where the diagonal Majorana phase matrix is given by  $P = \text{diag}(1, e^{i\alpha}, e^{i\beta})$ . Note that without this matrix the phase transformation required on the right side of  $U_{\text{PMNS}}$  (and hence on  $U_\nu$  itself) to bring it into the conventional phase structure of Eq. (3) is not possible, when one demands real positive diagonal neutrino mass entries in Eq. (6). The presence of  $P$  allows one to compensate for the phase transformation on the right side of  $U_{\text{PMNS}}$  without altering  $U_\nu$ . In this top-down approach one can then compare the  $U_{\text{PMNS}}$  obtained with  $U_{\text{TBM}}$ .

Instead, one might employ a bottom-up procedure to identify the  $\mu$ - $\tau$  symmetric neutrino mass matrix in the flavor basis as the most general one giving rise to tri-bimaximal mixing,

$$(m_\nu)_{\text{TBM}} = U_{\text{TBM}}^* P^* m_\nu^{\text{diag}} P^\dagger U_{\text{TBM}}^\dagger = \begin{pmatrix} A & B & B \\ \cdot & \frac{1}{2}(A+B+D) & \frac{1}{2}(A+B-D) \\ \cdot & \cdot & \frac{1}{2}(A+B+D) \end{pmatrix}. \quad (7)$$

The parameters  $A, B, D$  are in general complex and functions of the neutrino masses and Majorana phases:

$$A = \frac{1}{3} (2m_1 + m_2 e^{-2i\alpha}), \quad B = \frac{1}{3} (m_2 e^{-2i\alpha} - m_1), \quad D = m_3 e^{-2i\beta}. \quad (8)$$

Another way to write the mass matrix is to decompose it in terms of the three individual masses.

$$(m_\nu)_{\text{TBM}} = \frac{m_1}{6} \begin{pmatrix} 4 & -2 & -2 \\ \cdot & 1 & 1 \\ \cdot & \cdot & 1 \end{pmatrix} + \frac{m_2 e^{-2i\alpha}}{3} \begin{pmatrix} 1 & 1 & 1 \\ \cdot & 1 & 1 \\ \cdot & \cdot & 1 \end{pmatrix} + \frac{m_3 e^{-2i\beta}}{2} \begin{pmatrix} 0 & 0 & 0 \\ \cdot & 1 & -1 \\ \cdot & \cdot & 1 \end{pmatrix}. \quad (9)$$

Since observable departures from exact tri-bimaximal mixing are expected, it is of interest to study how deviations may arise. In a recent paper the authors linearly perturbed the matrix elements in Eq. (7) in order to observe how large the departures of the mixing angles from their tri-bimaximal values in Eq. (5) can be [13]. Allowing up to 20% deviations in the matrix elements in the normal hierarchy case results in non-zero values for  $\sin^2 \theta_{13}$  up to 0.001, for example. We argued that larger deviations in connection with a normal mass hierarchy would signal that the apparent nearly tri-bimaximal mixing should be considered accidental in nature, rather than the result of a softly-broken symmetry.

In this paper we study other mixing scenarios which deviate from tri-bimaximal mixing by leaving only one of the columns or one of the rows invariant. We shall refer to such scenarios as “generalized trimaximal mixing.” The term “trimaximal mixing” [14] was originally introduced to describe a mixing matrix in which only the second column of Eq. (4) remains invariant with the absolute value of every element of that column equal to  $1/\sqrt{3}$ . To distinguish the different versions, we label the original version of trimaximal mixing considered in [14–18] as  $\text{TM}_2$ , i.e.,

$$\text{TM}_2: \quad |U_{\alpha 2}|^2 = \frac{1}{3}, \quad \forall \alpha = e, \mu, \tau. \quad (10)$$

We shall also study the effects of allowing the other two columns to remain independently invariant under tri-bimaximal mixing perturbations and label them  $\text{TM}_1$  and  $\text{TM}_3$ , respectively. It is also of interest to allow one of the rows to remain invariant for which we adopt the labeling  $\text{TM}^i$ , with  $i = 1, 2$ , or  $3$ . In total there are six possibilities. By examining these

variations we shall learn that the three mixing angles and the CP-violating Dirac phase are not all independent and that restricted ranges of the deviations are imposed by the present mixing data. In fact, keeping a row or a column of  $U_{\text{PMNS}}$  fixed leads to a 2-parameter scenario, i.e., two of the four mixing observables in  $U_{\text{PMNS}}$  are determined. In four of the six trimaximal variants we find an interesting, characteristic and testable correlation between the mixing parameters. Future more precise data will allow one to test some of the trimaximal variations considered here. In particular, we stress that while  $\text{TM}_2$  allows naturally for non-zero  $\theta_{13}$ , non-maximal  $\theta_{23}$  and for  $\sin^2 \theta_{12} \neq \frac{1}{3}$ , it predicts that  $\sin^2 \theta_{12} \geq \frac{1}{3}$ , while the current best-fit points all lie at  $\sin^2 \theta_{12} \leq \frac{1}{3}$ . Some of the alternative trimaximal scenarios can accommodate this, while still allowing for non-zero  $\theta_{13}$ .

We would like to stress that in the present paper not merely deviations from tri-bimaximal mixing are discussed, but also novel and testable mixing scenarios are presented. Though tri-bimaximal mixing dominates the current phenomenological and theoretical discussion in neutrino physics, alternative proposals are surely of interest, and here we study one possible avenue.

In Sect. 2 we summarize the results obtained previously for broken tri-bimaximal mixing in [13]. In Sect. 3 we derive the results for the originally proposed trimaximal mixing and for the variant forms of trimaximal mixing in Sect. 4. A comparison of the results and conclusions are presented in Sect. 5.

## 2 Deviations from Tri-bimaximal Mixing

In [13] we raised the issue of deviations in the mixing observables from tri-bimaximal mixing when the mass matrix is perturbed. Our strategy was to perturb every entry of the mass matrix with a small complex parameter  $\epsilon_i$ , i.e:

$$m_\nu = \begin{pmatrix} A(1 + \epsilon_1) & B(1 + \epsilon_2) & B(1 + \epsilon_3) \\ \cdot & \frac{1}{2}(A + B + D)(1 + \epsilon_4) & \frac{1}{2}(A + B - D)(1 + \epsilon_5) \\ \cdot & \cdot & \frac{1}{2}(A + B + D)(1 + \epsilon_6) \end{pmatrix}, \quad (11)$$

where the complex perturbation parameters were taken to be  $|\epsilon_i| \leq 0.2$  for  $i = 1 - 6$  with their phases  $\phi_i$  allowed to lie between zero and  $2\pi$ . The results obtained for the oscillation parameters depend on the neutrino mass values and ordering. In case of a normal hierarchy ( $m_3 \gg m_2 > m_1$ ), the maximal expected values are

$$\text{NH:} \quad |U_{e3}| \lesssim \frac{2}{3} \sqrt{\frac{\Delta m_\odot^2}{\Delta m_A^2}} |\epsilon| \simeq 0.027, \quad \text{and} \quad \left| \frac{1}{2} - \sin^2 \theta_{23} \right| \lesssim \frac{1}{2} |\epsilon| \simeq 0.1, \quad (12)$$

where  $|\epsilon|$  denotes the absolute value of one of the six independent breaking parameters in Eq. (11). Fig. 1 shows for a maximal 20% deviation the parameter  $|U_{e3}|^2$  as a function of the smallest neutrino mass  $m_1$ . Note that the maximal value of  $|U_{e3}|$  depends linearly on  $\epsilon$ . Hence, if the deviation is allowed to range up to 50% for the normal hierarchy case, one finds that  $|U_{e3}|^2$  can reach a maximum value of roughly 0.005 for  $m_1 \lesssim 4$  meV. Even for a normal hierarchy it turns out that  $\sin^2 \theta_{12}$  can take values anywhere in its allowed range. The latter is also true in case of an inverted hierarchy ( $m_2 \simeq m_1 \gg m_3$ ), for which one furthermore

finds that the other mixing observables are bounded from above by

$$\text{IH:} \quad |U_{e3}| \lesssim \frac{1}{3}|\epsilon| \sqrt{\frac{8}{9} + \frac{16}{3} \frac{m_3}{\sqrt{\Delta m_A^2}}} \simeq 0.12, \quad \text{and} \quad \left| \frac{1}{2} - \sin^2 \theta_{23} \right| \lesssim \frac{8}{9}|\epsilon| \simeq 0.18. \quad (13)$$

Almost all of the  $3\sigma$  range can be covered. The interesting accessible range for  $|U_{e3}|^2$  as a function of the smallest neutrino mass  $m_3$  is also plotted in Fig. 1. If neutrinos are quasi-degenerate, then the fully allowed parameter space can be covered.

In Ref. [13] we also presented the results for some predictive  $SO(10)$  symmetric Grand Unified models exhibiting a normal mass hierarchy. There we found that  $|U_{e3}|^2 \gtrsim 2 \times 10^{-3}$ , while  $|\frac{1}{2} - \sin^2 \theta_{23}| \gtrsim 0.07$  with  $\sin^2 \theta_{12}$  typically  $< \frac{1}{3}$ . Hence one can distinguish the results of these models from softly-broken tri-bimaximal mixing once the value of  $|U_{e3}|^2$  and the neutrino mass hierarchy is known.

### 3 Trimaximal Mixing Deviations from TBM

We now turn to study the deviations from tri-bimaximal mixing which can arise with the less restrictive symmetry referred to as trimaximal mixing, defined by Eq. (10):

$$\text{TM}_2: \quad \begin{pmatrix} |U_{e2}|^2 \\ |U_{\mu 2}|^2 \\ |U_{\tau 2}|^2 \end{pmatrix} = \begin{pmatrix} 1/3 \\ 1/3 \\ 1/3 \end{pmatrix}. \quad (14)$$

This is the original version of trimaximality, for which also a model based on the  $\Delta(27)$  flavor symmetry was proposed [14, 18]. From the condition  $|U_{e2}|^2 = \frac{1}{3}$  it follows that

$$\sin^2 \theta_{12} = \frac{1}{3} \frac{1}{1 - |U_{e3}|^2} \geq \frac{1}{3}. \quad (15)$$

Note that  $\text{TM}_2$  predicts  $\sin^2 \theta_{12} \geq \frac{1}{3}$ , to be compared with the current best-fit values and  $1\sigma$  ranges of Ref. [8],  $\sin^2 \theta_{12} = 0.312_{-0.018}^{+0.019}$ , and Ref. [7]:  $\sin^2 \theta_{12} = 0.304_{-0.016}^{+0.022}$ . Inserting the range  $|U_{e3}|^2 = 0.016 \pm 0.010$  would give  $\sin^2 \theta_{12} = 0.339 \pm 0.003$ .

The second prediction of trimaximal mixing can be obtained from  $|U_{\mu 2}|^2 = \frac{1}{3}$  and is

$$\begin{aligned} \cos \delta \tan 2\theta_{23} &= \frac{2 \cos \theta_{13} \cot 2\theta_{13}}{\sqrt{2 - 3 \sin^2 \theta_{13}}} = \frac{1 - 2|U_{e3}|^2}{|U_{e3}| \sqrt{2 - 3|U_{e3}|^2}} \\ &\simeq \frac{1}{\sqrt{2}} \frac{1}{|U_{e3}|} \left( 1 - \frac{5}{4} |U_{e3}|^2 + \mathcal{O}(|U_{e3}|^4) \right). \end{aligned} \quad (16)$$

We observe that the three mixing angles are not independent but related as above. Because  $|U_{\tau 2}|^2$  is related by unitarity with  $|U_{e2}|^2$  and  $|U_{\mu 2}|^2$ , there is no third independent condition.

For  $\delta \neq \pi/2$  (or  $\delta \neq 3\pi/2$ ) and  $\theta_{13}$  non-zero,  $\theta_{23}$  is non-maximal. On the other hand, for maximal CP violation ( $\delta = \pi/2$  or  $3\pi/2$ ) it follows that  $\sin^2 \theta_{23} = \frac{1}{2}$ , independent of  $|U_{e3}|$ . We can use the expressions for  $\sin^2 \theta_{12}$  and  $\tan 2\theta_{23}$  to evaluate the Jarlskog invariant for leptonic CP violation, which in general reads  $J_{\text{CP}} = \text{Im}(U_{e1}^* U_{\mu 3}^* U_{e3} U_{\mu 1}) = \frac{1}{8} \sin 2\theta_{13} \cos \theta_{13} \sin 2\theta_{23} \sin 2\theta_{12} \sin \delta$ . We find

$$J_{\text{CP}} \simeq \frac{\sin 2\delta}{6\sqrt{2} \sqrt{\cos^2 \delta}} |U_{e3}|, \quad (17)$$

where we have expanded the lengthy exact equation, which is an odd function of  $\theta_{13}$ .

In Fig. 2 we show plots of  $\sin^2 \theta_{12}$  and  $\sin^2 \theta_{23}$  as functions of  $|U_{e3}|$  and  $\delta$ , respectively. As can be seen, solar neutrino mixing is well within the allowed  $2\sigma$  range, and the possible deviation from maximal atmospheric neutrino mixing is largest for CP conserving values of  $\delta$ , grows with  $|U_{e3}|$ , and can exceed the  $3\sigma$  range. Also displayed in Fig. 2 is a plot of  $\sin^2 \theta_{12}$  vs.  $\sin^2 \theta_{23}$ , when  $|U_{e3}|$  is allowed to vary. It is evident that the deviation from maximal atmospheric mixing can be larger than the deviation from  $\sin^2 \theta_{12} = \frac{1}{3}$ . Indeed, from  $\theta_{23} = \pi/4 - \epsilon$ , one obtains the leading order expressions  $\tan 2\theta_{23} \simeq \frac{1}{2\epsilon}$  and  $\sin^2 \theta_{23} \simeq \frac{1}{2} - \epsilon$ . By taking only the first order term on the RHS of Eq. (16) one obtains the relation

$$\left(\frac{1}{2} - \sin^2 \theta_{23}\right)^2 \simeq \frac{\cos^2 \delta}{2} |U_{e3}|^2 \simeq \frac{3 \cos^2 \delta}{2} \left(\sin^2 \theta_{12} - \frac{1}{3}\right). \quad (18)$$

This expression shows that the deviation from maximal atmospheric neutrino mixing can be stronger than the deviation from  $\sin^2 \theta_{12} = \frac{1}{3}$ . In general, the deviation in  $|U_{e3}|^2$  obtained here can be considerably larger than that entertained in Sect. 2.

The trimaximal mixing matrix can be parametrized by the application of a general 13-rotation from the right [17]

$$U_{\text{TM}_2} = U_{\text{TBM}} R_{13}(\theta; \psi), \quad \text{where } R_{13}(\theta; \psi) = \begin{pmatrix} \cos \theta & 0 & \sin \theta e^{-i\psi} \\ 0 & 1 & 0 \\ -\sin \theta e^{i\psi} & 0 & \cos \theta \end{pmatrix}. \quad (19)$$

It is easy to see that  $|U_{e3}|^2 = \frac{2}{3} \sin^2 \theta$  and  $\sin^2 \theta_{12} = 1/(3 - 2 \sin^2 \theta)$ , which is equal to  $\frac{1}{3}/(1 - |U_{e3}|^2)$  as before. We find furthermore that

$$\sin^2 \theta_{23} = \frac{1}{2} + \frac{1}{2\sqrt{3}} \frac{\sin 2\theta \cos \psi}{1 - |U_{e3}|^2} \quad \text{and} \quad J_{\text{CP}} = \frac{\sin 2\theta \sin \psi}{6\sqrt{3}}, \quad (20)$$

where  $\sin 2\theta = \sqrt{6} |U_{e3}| \sqrt{1 - \frac{3}{2} |U_{e3}|^2}$ . The CP phase  $\delta$  in this parameterization is related to the phase  $\psi$ .

By using the above form for  $U_{\text{TM}_2}$  in Eq. (19), we can obtain the corresponding neutrino mass matrix in the lepton flavor basis from

$$(m_\nu)_{\text{TM}_2} = U_{\text{TM}_2}^* P^* m_\nu^{\text{diag}} P^\dagger U_{\text{TM}_2}^\dagger = \begin{pmatrix} A & B+C & B-C \\ \cdot & \frac{1}{2}(A+B+D-2C) & \frac{1}{2}(A+B-D) \\ \cdot & \cdot & \frac{1}{2}(A+B+D+2C) \end{pmatrix}$$

where we identify (with  $c_\theta = \cos \theta$  and  $s_\theta = \sin \theta$ )

$$\begin{aligned} A &= \frac{1}{3} \left( 2m_1 c_\theta^2 + 2m_3 s_\theta^2 e^{2i(\psi-\beta)} + m_2 e^{-2i\alpha} \right), & C &= \frac{1}{\sqrt{3}} \left( m_1 e^{-i\psi} - m_3 e^{i(\psi-2\beta)} \right) s_\theta c_\theta, \\ B &= \frac{1}{3} \left( m_2 e^{-2i\alpha} - m_1 c_\theta^2 - m_3 s_\theta^2 e^{2i(\psi-\beta)} \right), & D &= m_3 e^{-2i\beta} c_\theta^2 + m_1 e^{-2i\psi} s_\theta^2. \end{aligned}$$

It is clear from this mass matrix that the additional  $C$  terms break the original  $\mu$ - $\tau$  symmetry present with tri-bimaximal mixing in a well-defined way. Note that  $C$  vanishes for  $\theta = 0$  and that in this case  $(m_\nu)_{\text{TBM}}$  from Eq. (7) is recovered. Interestingly, if we decompose the mass matrix for  $\text{TM}_2$  in terms of the individual neutrino masses, as done for TBM in Eq. (9), we find that  $m_2$  is multiplied with the same flavor-democratic matrix as in Eq. (9). The other

two masses are multiplied now with more complicated matrices, having entries depending on the angle  $\theta$ .

Since we can now trade  $\theta_{12}$  for  $\theta_{13}$  with the use of Eq. (15), it is possible to obtain a simple value for the effective mass  $\langle m_{ee} \rangle$  governing neutrinoless double beta decay ( $0\nu\beta\beta$ ). With inverted hierarchy ( $m_2 \simeq m_1 \gg m_3$ ) the general result is  $\langle m \rangle = c_{13}^2 \sqrt{\Delta m_A^2} \sqrt{1 - \sin^2 2\theta_{12} \sin^2 \alpha}$ . In case of  $\text{TM}_2$  mixing we then find

$$\langle m \rangle \simeq \sqrt{\Delta m_A^2} \left[ \sqrt{1 - \frac{8}{9} \sin^2 \alpha} - \frac{1 - \frac{2}{3} \sin^2 \alpha}{\sqrt{1 - \frac{8}{9} \sin^2 \alpha}} |U_{e3}|^2 \right]. \quad (21)$$

The  $\frac{8}{9}$  in the first term is of course the value of  $\sin^2 2\theta_{12}$  in case of tri-bimaximal mixing.

## 4 Variant Trimaximal Mixing Scenarios

We have seen that the prediction of trimaximality implies  $\sin^2 \theta_{12} \geq \frac{1}{3}$ , while the best-fit points are all below  $\frac{1}{3}$ . If this trend continues, then one may introduce variants of trimaximal mixing, such as

$$\text{TM}_1 : \quad \begin{pmatrix} |U_{e1}|^2 \\ |U_{\mu 1}|^2 \\ |U_{\tau 1}|^2 \end{pmatrix} = \begin{pmatrix} 2/3 \\ 1/6 \\ 1/6 \end{pmatrix}. \quad (22)$$

Here we have fixed the first column of  $U_{\text{PMNS}}$  to have the same form as in the case of tri-bimaximal mixing<sup>2</sup>. A consequence of Eq. (22) is that  $\sin^2 \theta_{12} \leq \frac{1}{3}$ . Indeed, from  $|U_{e1}|^2 = \frac{2}{3}$  one finds

$$\sin^2 \theta_{12} = \frac{1}{3} \frac{1 - 3|U_{e3}|^2}{1 - |U_{e3}|^2} \simeq \frac{1}{3} (1 - 2|U_{e3}|^2). \quad (23)$$

Using the range  $|U_{e3}|^2 = 0.016 \pm 0.010$  gives  $\sin^2 \theta_{12} = 0.322 \pm 0.007$ . Fig. 3 shows  $\sin^2 \theta_{12}$  as a function of  $|U_{e3}|$ . In contrast to the original trimaximal mixing scheme,  $\text{TM}_2$ , the best-fit value of  $\sin^2 \theta_{12}$  can be obtained (for  $|U_{e3}| \simeq 0.179$ ). Note that  $\sin^2 \theta_{12}$  decreases with  $|U_{e3}|$ . As long as  $|U_{e3}| \gtrsim 0.06$ ,  $\sin^2 \theta_{12}$  is within its allowed  $1\sigma$  range.

The second independent condition in Eq. (22) involving  $|U_{\mu 1}|^2 = 1/6$  gives

$$\cos \delta \tan 2\theta_{23} = -\frac{1 - 5|U_{e3}|^2}{2\sqrt{2}|U_{e3}|\sqrt{1 - 3|U_{e3}|^2}} \simeq \frac{-1}{2\sqrt{2}|U_{e3}|} \left( 1 - \frac{7}{2}|U_{e3}|^2 \right). \quad (24)$$

Here the results are qualitatively similar to the ones for trimaximal mixing treated in Sect. 3. In Fig. 3 we also show  $\sin^2 \theta_{23}$  as a function of  $\delta$ . As for  $\text{TM}_2$  mixing, CP conserving values of  $\delta$  maximize the deviations, but here the  $3\sigma$  range can easily be overshoot. Finally, in Fig. 3,  $\sin^2 \theta_{23}$  is plotted against  $\sin^2 \theta_{12}$ . Again, the possible departure from  $\sin^2 \theta_{23} = \frac{1}{2}$  can be larger than the one from  $\sin^2 \theta_{12} = \frac{1}{3}$ .

A mixing matrix with the  $\text{TM}_1$  property can be obtained by multiplying  $U_{\text{TBM}}$  with a 23-rotation of angle  $\theta$  from the right,  $U_{\text{TM}_1} = U_{\text{TBM}} R_{23}(\theta; \psi)$ . The observables are in this case

$$\begin{aligned} |U_{e3}|^2 &= \frac{1}{3} \sin^2 \theta, \quad \sin^2 \theta_{23} = \frac{1}{2} - \frac{\sqrt{\frac{3}{2}} \sin 2\theta \cos \psi}{3 - \sin^2 \theta}, \\ \sin^2 \theta_{12} &= 1 - \frac{2}{3 - \sin^2 \theta}, \quad J_{\text{CP}} = \frac{1}{6\sqrt{6}} \sin 2\theta \sin \psi, \end{aligned} \quad (25)$$

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<sup>2</sup>This possibility has been mentioned first in Ref. [19], but its phenomenology has not been studied yet.

where the CP phase  $\delta$  is again related to  $\psi$ . For the mass matrix we find

$$(m_\nu)_{\text{TM}_1} = U_{\text{TM}_1}^* P^* m_\nu^{\text{diag}} P^\dagger U_{\text{TM}_1}^\dagger = \begin{pmatrix} A & B+C & B-C \\ \cdot & \frac{1}{2}(A+B+D+4C) & \frac{1}{2}(A+B-D) \\ \cdot & \cdot & \frac{1}{2}(A+B+D-4C) \end{pmatrix},$$

where we identify

$$\begin{aligned} A &= \frac{1}{3} \left( 2m_1 + m_2 c_\theta^2 e^{-2i\alpha} + m_3 s_\theta^2 e^{2i(\psi-\beta)} \right), \quad C = \frac{1}{\sqrt{6}} \left( m_2 e^{-i(\psi+2\alpha)} - m_3 e^{i(\psi-2\beta)} \right) s_\theta c_\theta, \\ B &= \frac{1}{3} \left( -m_1 + m_2 c_\theta^2 e^{-2i\alpha} + m_3 s_\theta^2 e^{2i(\psi-\beta)} \right), \quad D = m_2 e^{-2i(\psi+\alpha)} s_\theta^2 + m_3 c_\theta^2 e^{-2i\beta}. \end{aligned}$$

Again we see that the original  $\mu$ - $\tau$  symmetry is broken by the extra terms involving  $C$ . We can decompose the mass matrix for  $\text{TM}_1$  in terms of the individual neutrino masses and find that  $m_1$  is multiplied with the same matrix as in Eq. (9).

If we would insist that the third column of  $U_{\text{TBM}}$  remain invariant instead, i.e.,  $|U_{e3}|^2 = 0$ ,  $|U_{\mu 3}|^2 = |U_{\tau 3}|^2 = \frac{1}{2}$ , then  $\theta_{13} = \delta = 0$ ,  $\theta_{23} = \pi/4$ , while  $\theta_{12}$  is a free parameter. This case ( $\text{TM}_3$  in our notation) is nothing other than the well-known  $\mu$ - $\tau$  symmetry.

It is also of interest to consider the case where one of the rows of the tri-bimaximal mixing matrix remains invariant. Again such a result can be obtained by multiplying with a suitable two dimensional rotation matrix, but now from the left. In fact, this class of deviations from tri-bimaximal mixing corresponds to the charged lepton flavor matrix differing from its diagonal mass matrix, thus introducing a  $U_\ell^\dagger$  factor in the PMNS mixing matrix, i.e.,  $U_{\text{PMNS}} = U_\ell^\dagger U_{\text{TBM}}$  [20].

Let us start with the case of the first row in  $U_{\text{TBM}}$  remaining invariant. We denote this with a superscript as

$$\text{TM}^1 : \quad \left( |U_{e1}|^2, |U_{e2}|^2, |U_{e3}|^2 \right) = \left( \frac{2}{3}, \frac{1}{3}, 0 \right), \quad (26)$$

As a result,  $\theta_{23}$  is a free parameter, while  $\sin^2 \theta_{12} = \frac{1}{3}$ , as well as  $\theta_{13} = \delta = 0$ . With a rotation of  $U_{\text{TBM}}$  by the matrix  $R_{23}(\theta; \psi)$  from the left, the light Majorana neutrino mass matrix becomes

$$(m_\nu)_{\text{TM}^1} = R_{23}^* (m_\nu)_{\text{TBM}} R_{23}^\dagger. \quad (27)$$

If we consider the second or third row we can correlate all four mixing parameters. Starting with the second row, i.e.,

$$\text{TM}^2 : \quad \left( |U_{\mu 1}|^2, |U_{\mu 2}|^2, |U_{\mu 3}|^2 \right) = \left( \frac{1}{6}, \frac{1}{3}, \frac{1}{2} \right), \quad (28)$$

one immediately finds from  $|U_{\mu 3}|^2 = \frac{1}{2}$ :

$$\sin^2 \theta_{23} = \frac{1}{2(1 - |U_{e3}|^2)} \simeq \frac{1}{2} (1 + |U_{e3}|^2) \geq \frac{1}{2}, \quad (29)$$

i.e., atmospheric neutrino mixing on the “dark side,” with a maximal value of  $\sin^2 \theta_{23} \simeq 0.524$  for  $|U_{e3}|^2 = 0.046$ . Inserting Eq. (29) in  $|U_{\mu 2}|^2 = \frac{1}{3}$  gives a complicated and lengthy expression including  $\cos \delta$ ,  $\sin^2 \theta_{12}$  and  $|U_{e3}|$ , which can be approximated as

$$\sin^2 \theta_{12} \simeq \frac{1}{3} - \frac{2\sqrt{2}}{3} |U_{e3}| \cos \delta + \frac{1}{3} |U_{e3}|^2 \cos 2\delta. \quad (30)$$



Fig. 4 shows the result for the  $\text{TM}^2$  scenario. The parameter dependence is similar to the one for the  $\text{TM}_1$  and  $\text{TM}_2$  scenarios, the main difference being the exchanged roles of  $\theta_{12}$  and  $\theta_{23}$ .

On the other hand with the third row remaining invariant under the transformation

$$\text{TM}^3: \left( |U_{\tau 1}|^2, |U_{\tau 2}|^2, |U_{\tau 3}|^2 \right) = \left( \frac{1}{6}, \frac{1}{3}, \frac{1}{2} \right) \quad (31)$$

the same approximate formula with a relative sign for the term of order  $|U_{e3}|$  is found. Now, however, one finds atmospheric neutrino mixing on the “bright side”:

$$\sin^2 \theta_{23} = \frac{1 - 2|U_{e3}|^2}{2(1 - |U_{e3}|^2)} \simeq \frac{1}{2} (1 - |U_{e3}|^2) \leq \frac{1}{2}. \quad (32)$$

The maximal deviation occurs for the largest possible  $|U_{e3}|^2 = 0.046$ , in which case  $\sin^2 \theta_{23} \simeq 0.476$ . In Fig. 5 the resulting correlations for the  $\text{TM}^3$  scenario are given.

As in the case of  $\text{TM}^1$ , the  $U_{\text{TM}^2}$  mixing matrix can be obtained by a rotation of  $U_{\text{TBM}}$  from the left by  $R_{13}$ . For  $\text{TM}^3$ , the rotation matrix is  $R_{12}$ . For all three cases with the first, second, or third row of  $U_{\text{TBM}}$  remaining invariant, solar neutrino mixing can occur in the complete  $3\sigma$  range. The deviations from tri-bimaximal mixing can be considerably larger for  $\theta_{13}$  in the normal hierarchical case than were observed in Sect. 2 and Ref. [13].

Finally, it is amusing to compare the trimaximal mixing schemes with the recently proposed tetramaximal one [21]. Its name stems from the fact that it can be obtained by four consecutive rotations each having a maximal angle of  $\pi/4$ , and with appropriately chosen phases:  $U_{\text{tetra}} = R_{23}(\pi/4; \pi/2) R_{13}(\pi/4; 0) R_{12}(\pi/4; 0) R_{13}(\pi/4; \pi)$ . The predictions are  $\delta = \pi/2$ ,  $\sin^2 \theta_{23} = \frac{1}{2}$ ,  $|U_{e3}|^2 = \frac{1}{4}(\frac{3}{2} - \sqrt{2}) \simeq 0.021$  and  $\sin^2 \theta_{12} = 1/(\frac{5}{2} + \sqrt{2}) \simeq 0.255$ . None of the trimaximal variants discussed in this paper could be confused with tetramaximal mixing.

## 5 Summary and Conclusions

Although the present PMNS lepton mixing matrix deduced from experiment is consistent with tri-bimaximal mixing, it is of interest to study possible deviations which may arise in the future. In a previous paper the authors considered linear complex perturbations of the neutrino matrix away from the most general  $\mu$ - $\tau$  symmetric texture which yields tri-bimaximal mixing. Here we have considered variations of trimaximal mixing which can arise, e.g., with a simple complex rotation of the  $U_{\text{TBM}}$  matrix from the right or the left. The original trimaximal mixing matrix  $U_{\text{TM}^2}$  preserves the second column of  $U_{\text{TBM}}$  for which each element has absolute value of  $1/\sqrt{3}$ . We have generalized this mixing scenario here and refer to the other variations as  $\text{TM}_k$  or  $\text{TM}^k$  according to which  $k = 1, 2, 3$  column or row remains invariant, respectively. Independent of our comparisons with tri-bimaximal mixing, the “trimaximal” mixing scenarios considered here are alternative, novel and testable mixing schemes.

We summarize our findings in Table 1 for the squares of the sines of the three mixing angles. For broken TBM where up to 20% deviations are allowed in every neutrino mass matrix element, we had found that the allowed variations in sine squared of the mixing angles are uncorrelated to a large extent. The perturbed solar and atmospheric neutrino mixing angles cover the entire mixing ranges presently allowed, while  $|U_{e3}|^2 = \sin^2 \theta_{13}$  can range from zero up to 0.001 (0.014) for the case of normal (inverted) hierarchy. These upper

	$ U_{e3} ^2$	$\sin^2 \theta_{12}$	$\sin^2 \theta_{23}$
TBM	0	$\frac{1}{3}$	$\frac{1}{2}$
Broken TBM			
NH	$\lesssim 0.001$	$0.26 - 0.38$	$0.37 - 0.63$
IH	$\lesssim 0.014$	$0.26 - 0.38$	$0.35 - 0.64$
TM <sub>1</sub>	$\leq 0.046$	$\frac{1}{3}(1 - 2 U_{e3} ^2)$ $0.30 - 0.33$	$\frac{1}{2} - \sqrt{2} U_{e3} (1 - \frac{1}{2} U_{e3} ^2) \cos \delta$ $0.34 - 0.64$
TM <sub>2</sub>	$\leq 0.046$	$\frac{1}{3}(1 +  U_{e3} ^2)$ $0.33 - 0.35$	$\frac{1}{2} + \frac{1}{\sqrt{2}} U_{e3} (1 + \frac{1}{4} U_{e3} ^2) \cos \delta$ $0.34 - 0.64$
TM <sub>3</sub>	0	$0.26 - 0.38$	$\frac{1}{2}$
TM <sup>1</sup>	0	$\frac{1}{3}$	$0.34 - 0.64$
TM <sup>2</sup>	$\leq 0.046$	$\frac{1}{3} - \frac{2\sqrt{2}}{3} U_{e3}  \cos \delta + \frac{1}{3} U_{e3} ^2 \cos 2\delta$ $0.26 - 0.38$	$\frac{1}{2}(1 +  U_{e3} ^2)$ $0.50 - 0.52$
TM <sup>3</sup>	$\leq 0.046$	$\frac{1}{3} + \frac{2\sqrt{2}}{3} U_{e3}  \cos \delta + \frac{1}{3} U_{e3} ^2 \cos 2\delta$ $0.26 - 0.38$	$\frac{1}{2}(1 -  U_{e3} ^2)$ $0.49 - 0.50$

Table 1: Summary of the mixing angles obtained in the exact and broken tri-bimaximal, and generalized trimaximal mixing schemes. We have for the sake of illustration expanded the exact correlations (see Figs. 2 – 5). For the ranges given, the present  $3\sigma$  data bounds on the three mixing angles have been imposed.

bounds depend sensitively on the lightest neutrino mass as shown in Fig. 1, especially for a normal ordering. The largest upper bounds presently allowed are reached in the case of three-fold neutrino mass degeneracy.

For the six generalized trimaximal mixing cases considered, on the other hand, the mixing angles and Dirac CP phase are characteristically correlated for four of the cases. The exceptional cases arise for TM<sub>3</sub> and TM<sup>1</sup>. For TM<sub>3</sub> both  $|U_{e3}|^2$  and  $\sin^2 \theta_{23}$  remain fixed at their TBM values, while  $\theta_{12}$  is a free variable and limited only by the present experimental bounds. This case corresponds to the well-known  $\mu$ - $\tau$  symmetry. For TM<sup>1</sup>,  $|U_{e3}|^2$  and  $\sin^2 \theta_{12}$  remain fixed, while  $\sin^2 \theta_{23}$  remains bounded only by experiment. For TM<sub>1</sub> and TM<sub>2</sub> the solar mixing is tightly limited, with the former ranging just below the TBM value and the latter just above the TBM value of  $\frac{1}{3}$ . The atmospheric neutrino mixing in these two cases can cover the full presently allowed region with the former (latter) peaking at  $\delta = 0$  ( $\pi$ ) and bottoming at  $\delta = \pi$  (0). For TM<sup>2</sup> and TM<sup>3</sup> the opposite situation holds, where the atmospheric neutrino mixing range is tightly limited close to the TBM value in the bright side or dark side, respectively, while the full range for solar neutrino mixing can be realized.

When more refined ranges for the mixing angles are known, one will be able to rule out or confirm the new mixing scenarios discussed here and narrow down the acceptable deviations

from tri-bimaximal mixing we have discussed in this paper.

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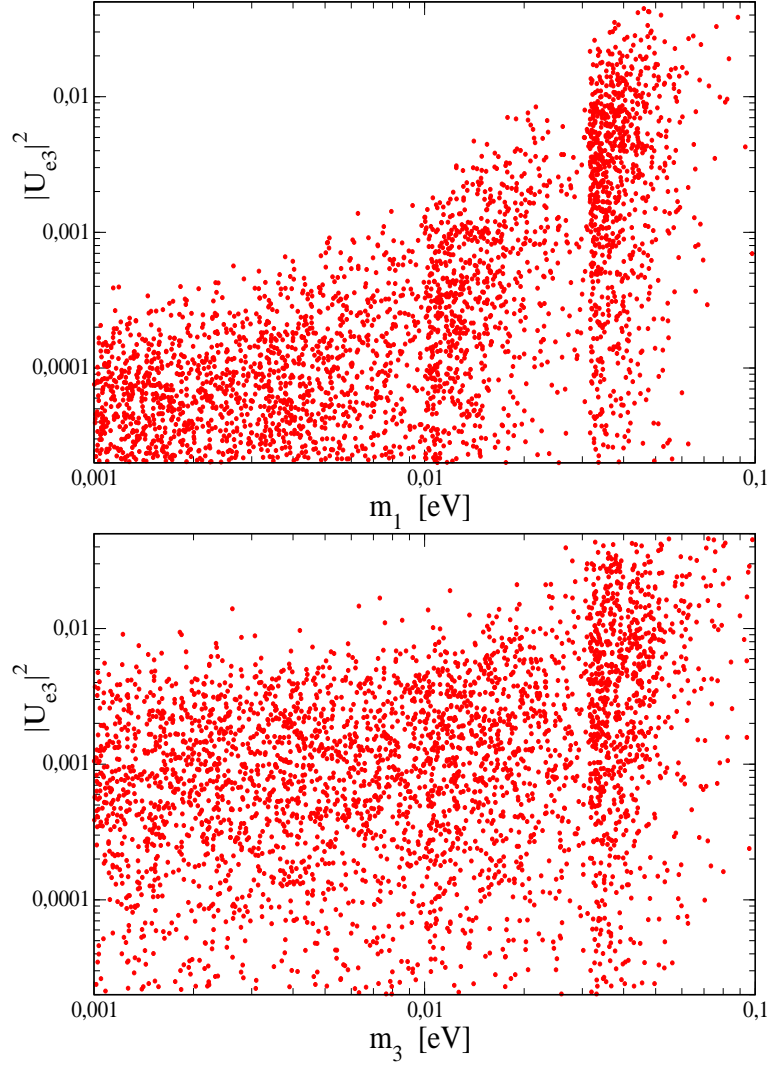


Figure 1: Scatter plot of the predictions for  $|U_{e3}|^2$  versus the smallest neutrino mass for softly broken tri-bimaximal mixing. Shown are the normal mass ordering (upper plot) and the inverted one (lower plot).

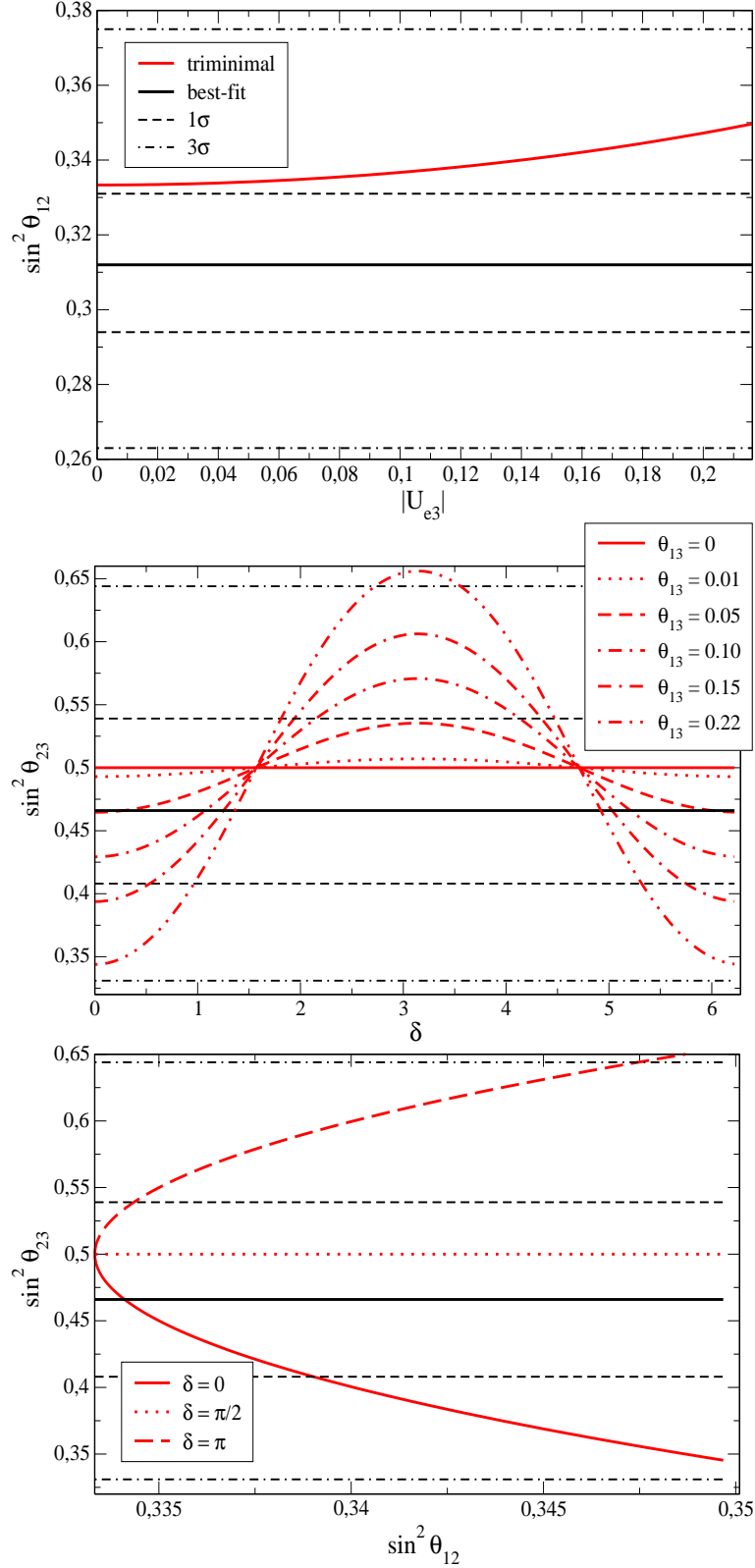


Figure 2: Phenomenology of exact  $TM_2$  mixing. Shown are the solar neutrino parameter  $\sin^2 \theta_{12}$  against  $|U_{e3}|$ , the atmospheric neutrino parameter  $\sin^2 \theta_{23}$  against  $\delta$  for different values of  $|U_{e3}|$  and  $\sin^2 \theta_{12}$  against  $\sin^2 \theta_{23}$ . Also given are the current best-fit value and the  $1\sigma$  as well as  $3\sigma$  ranges from a global fit [8].

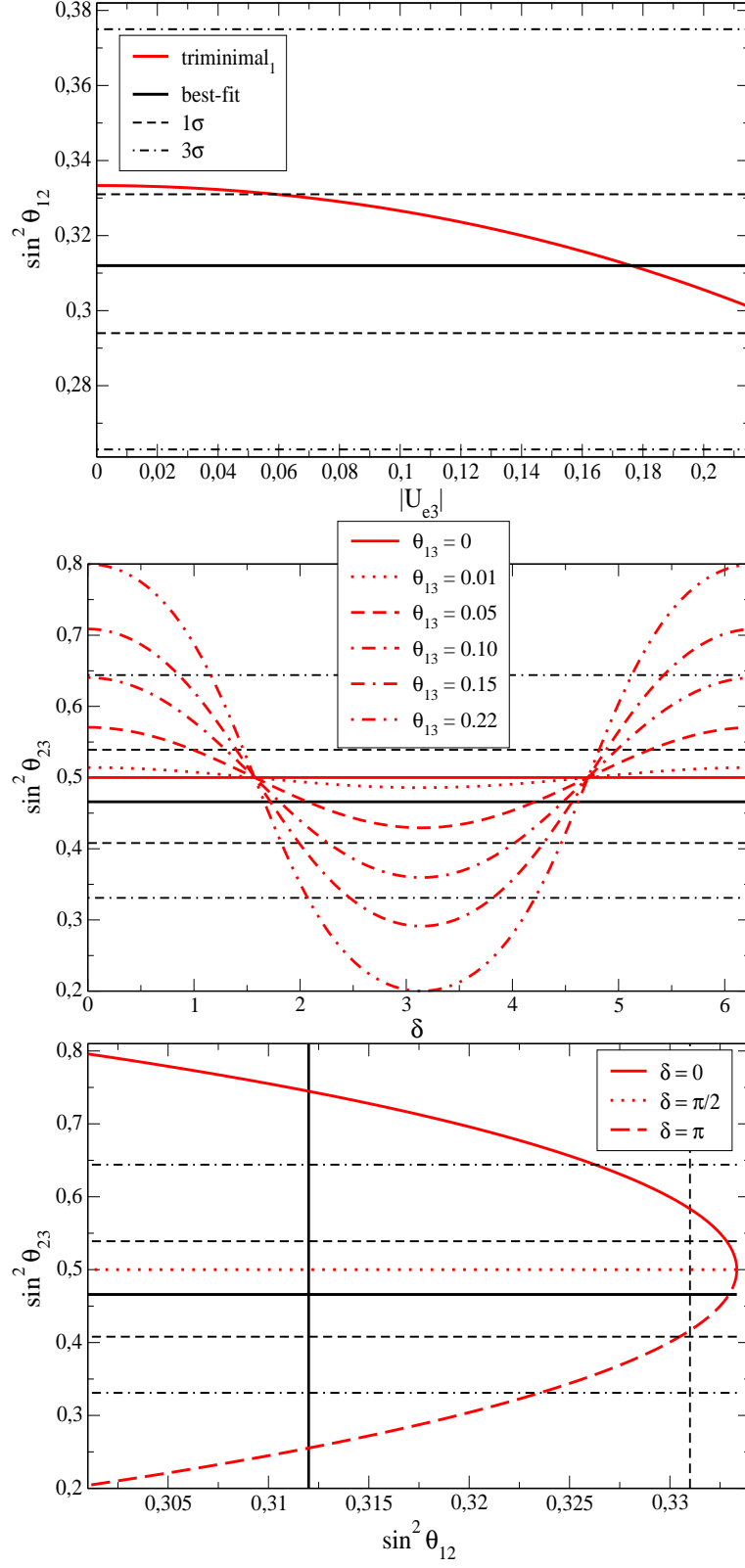


Figure 3: Phenomenology of exact  $TM_1$  mixing. Shown are the solar neutrino parameter  $\sin^2 \theta_{12}$  against  $|U_{e3}|$ , the atmospheric neutrino parameter  $\sin^2 \theta_{23}$  against  $\delta$  for different values of  $|U_{e3}|$  and  $\sin^2 \theta_{12}$  against  $\sin^2 \theta_{23}$ . Also given are the current best-fit value and the  $1\sigma$  as well as  $3\sigma$  ranges from a global fit [8].

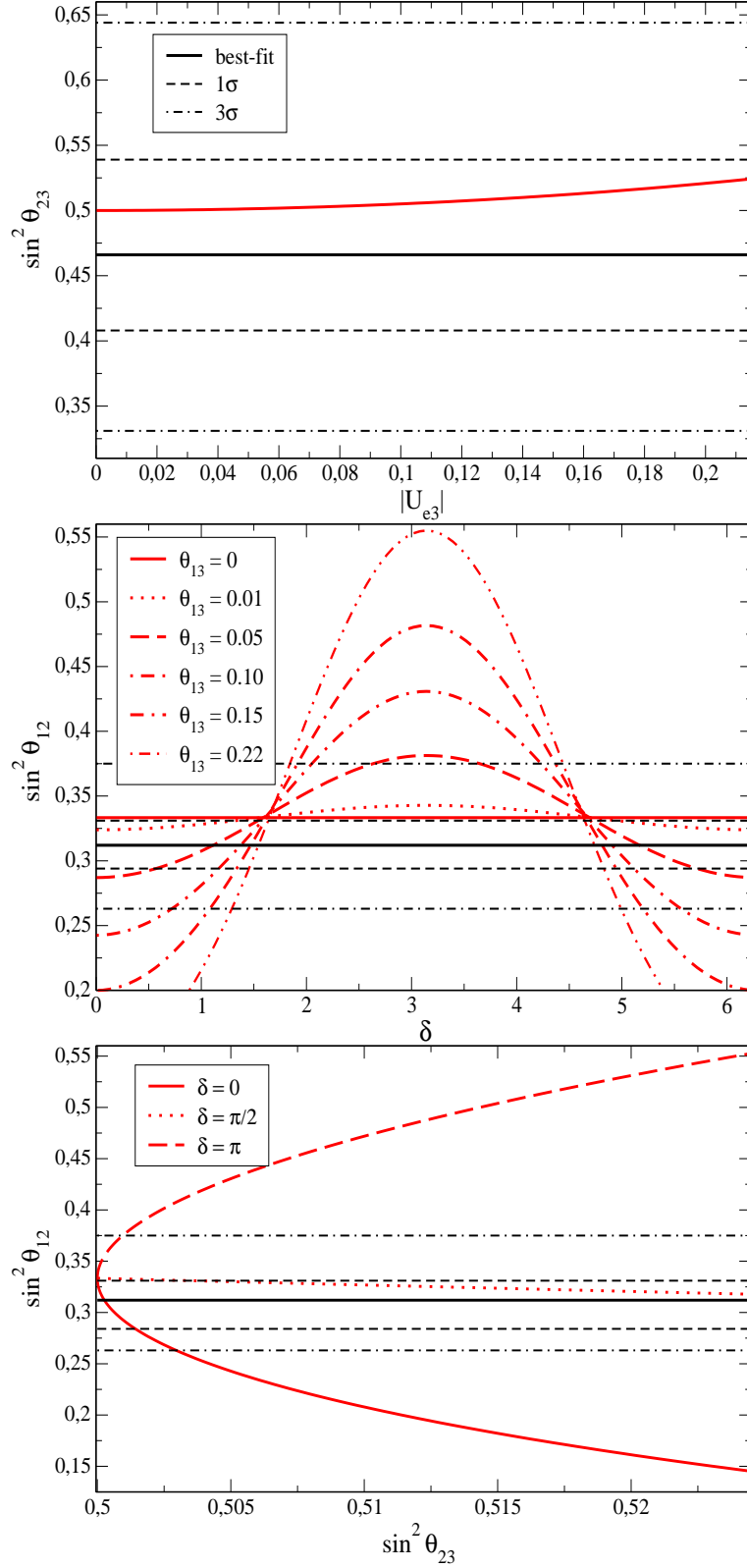


Figure 4: Phenomenology of exact  $TM^2$  mixing. Shown are the atmospheric neutrino parameter  $\sin^2 \theta_{23}$  against  $|U_{e3}|$ , the solar neutrino parameter  $\sin^2 \theta_{12}$  against  $\delta$  for different values of  $|U_{e3}|$  and  $\sin^2 \theta_{12}$  against  $\sin^2 \theta_{23}$ . Also given are the current best-fit value and the  $1\sigma$  as well as  $3\sigma$  ranges from global fit [8].



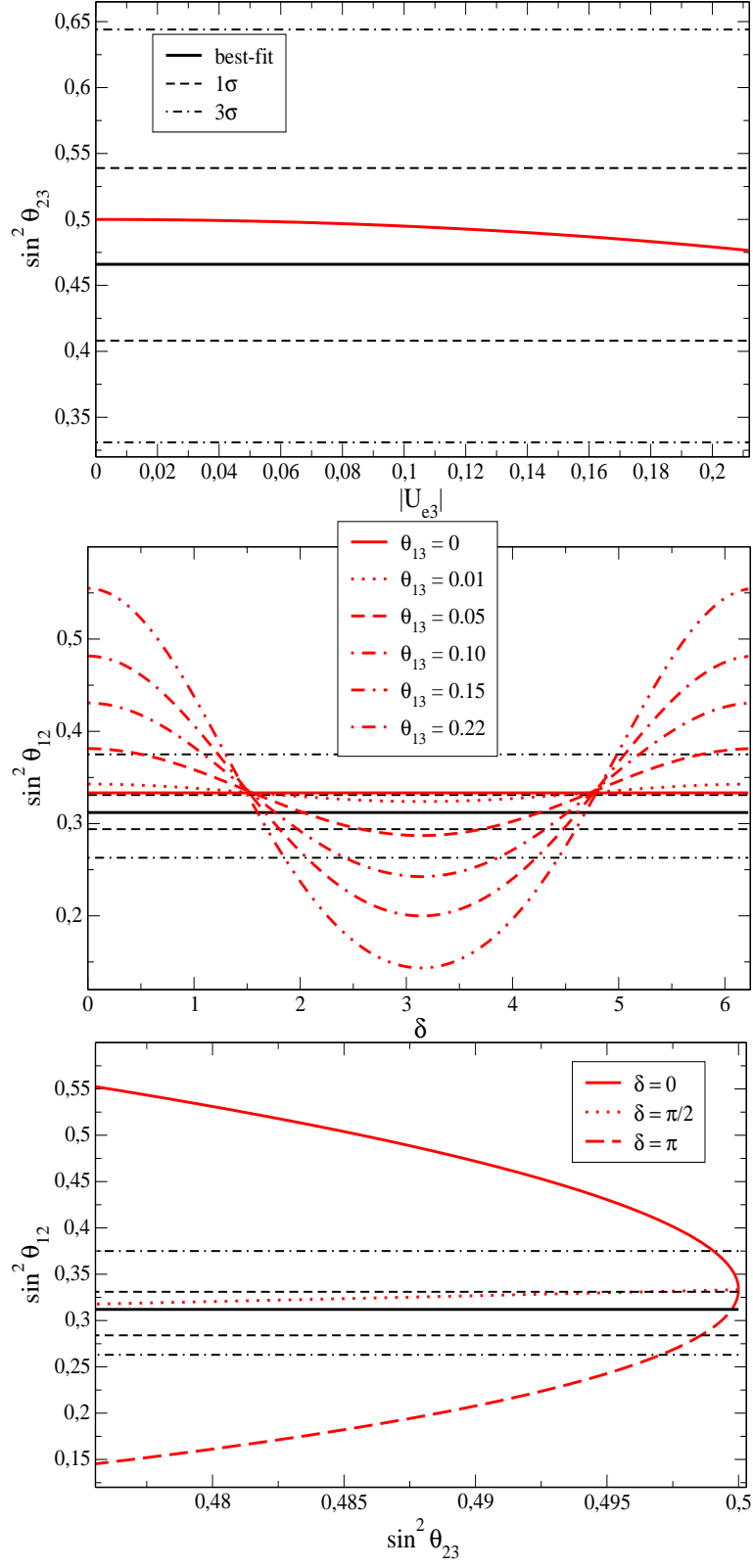


Figure 5: Phenomenology of exact  $TM^3$  mixing. Shown are the atmospheric neutrino parameter  $\sin^2 \theta_{23}$  against  $|U_{e3}|$ , the solar neutrino parameter  $\sin^2 \theta_{12}$  against  $\delta$  for different values of  $|U_{e3}|$  and  $\sin^2 \theta_{12}$  against  $\sin^2 \theta_{23}$ . Also given are the current best-fit value and the  $1\sigma$  as well as  $3\sigma$  ranges from a global fit [8].